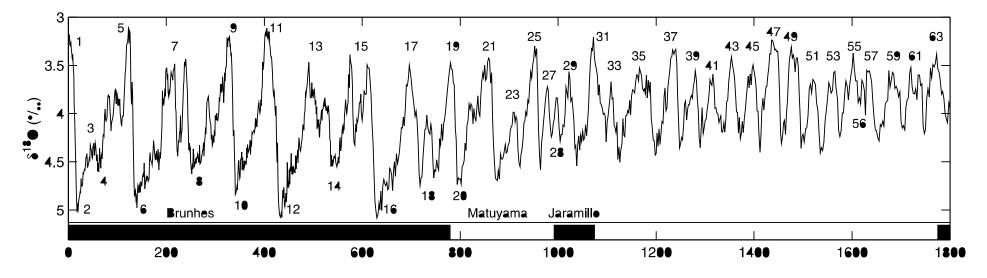
# Meta-stablity in stochastic partial differential equations induced by Lévy noise

Claudia Hein Michael Högele Peter Imkeller Ilya Pavlyukevich

Humboldt-Universität zu Berlin Institut für Mathematik

Lille, September 1, 2008



#### **1.** Paleoclimatic time series

Lisiecki, Raymo, Paleoceanography 2005 concentration variation of <sup>18</sup>O to <sup>16</sup>O taken from marine sediments at 57 globally distributed sites (e.g. Brunhes, Matuyama, Jaramillo):

global average temperature time series

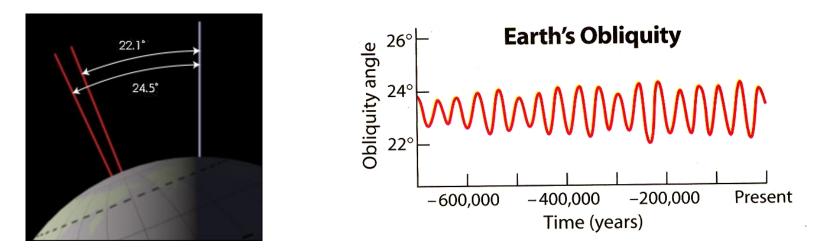
basic feature:

- from 0 to -1 Myr periodicity  $\sim$  100 000 y
- from -1 Myr to -1.8 Myr periodicity  $\sim$  44 000 y

# 2. Milankovich cycles

Milankovich (1920): astronomical perturbations of earth's trajectory give rise to basically three cycles: precession, obliquity, eccentricity

#### obliquity (axial tilt) (41 000 y)

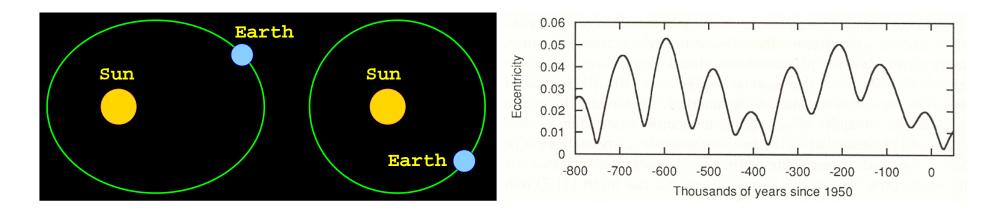


periodic wobbling of Earth's rotation axis: 2.4 deg change in tilt of axis with respect to plane of orbit; periodicity: 41 000 y; present tilt: 23.4 deg.

increase of obliquity: increase of amplitude of seasonal cycle in insolation, summers receiving more radiative flux, winters less

## 3. Milankovich cycles

#### eccentricity (100 000 y)



eccentricity: measure of deviation of ellipse from circularity; periodic variation between 0.005 and 0.058; reason: interactions with gravitational fields of Jupiter and Saturn; periodicity 100,000 y; present eccentricity: 0.017.

Orbital mechanics: eccentricity extreme means seasons on far side of orbit last substantially longer.

# 4. EBM: Gaussian SDE model; noise induced transitions

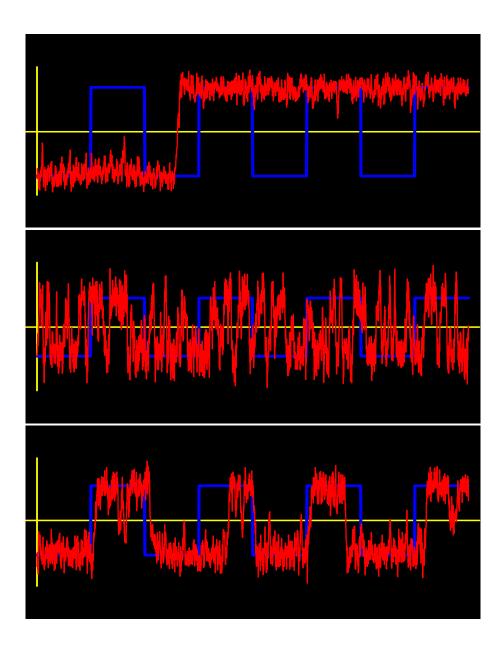
C. Nicolis; Benzi et al. 80' Energy Balance Model for global Earth temperature X: attempt of simple explanation of time series

$$c\frac{dX(t)}{dt} = R_{\rm in} - R_{\rm out}$$

Gaussian SDE

$$dX(t) = -U'\left(\frac{t}{T}, X(t)\right) dt + \varepsilon \, dW(t)$$

Double-well potential with periodically varying wells' depths  $U'(t, \pm 1) = U'(t, 0) = 0$ Period T – large; noise  $\varepsilon$  – small. Different noise intensities lead to different regimes.



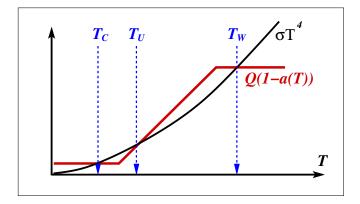
# 5. Gaussian transition times

#### simplest model

slowly time varying potential gradient

$$U'(t,x) = Q_t(1-a(x)) - \sigma x^4,$$

first omit time variation; double well potential U;



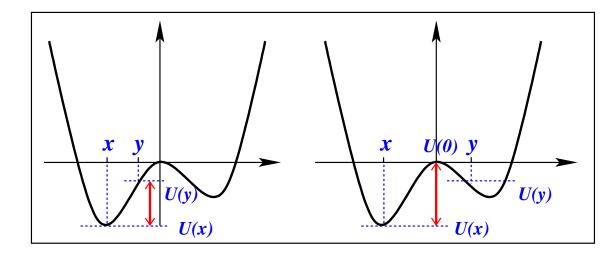
consider sde

$$dX^{\varepsilon}(t) = -U'(X^{\varepsilon}(t))dt + \varepsilon dW(t)$$

Under which conditions transitions can be expected? How long do they take in the small noise limit?

#### 6. Gaussian transition times

action functional  $\Rightarrow$  pseudopotential V(x, y) = 2[work  $x \rightarrow y],$ 



 $V(x,y) = 2[U(y) - U(x)]^+ \quad V(x,y) = 2[U(0) - U(x)]^+$  $\tau_y^{\varepsilon} = \inf\{t \ge 0 : X^{\varepsilon}(t) = y\} \quad \text{transition time}$ 

**Thm 1**(Freidlin, transition law) for all  $\delta$ 

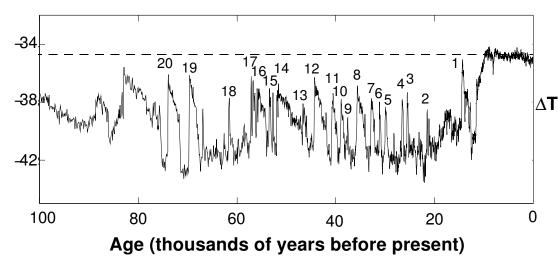
$$\mathbf{P}_{x}\left[\exp\left(\frac{V(x,y)-\delta}{\varepsilon^{2}}\right) \leq \tau_{y}^{\varepsilon} \leq \exp\left(\frac{V(x,y)+\delta}{\varepsilon^{2}}\right)\right] \rightarrow_{\varepsilon \to 0} 1$$

interprets Kramers-Eyring law:  $E_x(\tau_y^{\varepsilon}) \sim \exp(\frac{V(x,y)}{\varepsilon^2})$ 

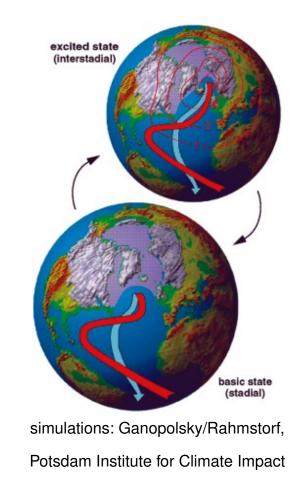
## 7. Dansgaard-Oeschger events

temperature indicators: <sup>18</sup>O, <sup>16</sup>O, methane, calcium etc.

GRIP ice core data: 20 abrupt changes in climate of Greenland during last ice age (-91 000 to -11 000 y) (D/O events).



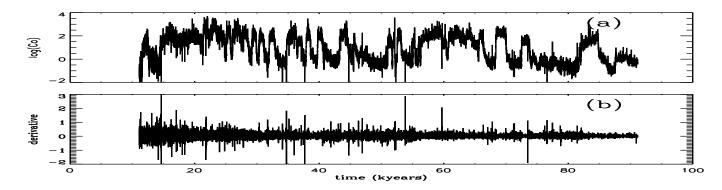
- rapid warming by 5-10°C within one decade
- subsequent slower cooling within a few centuries
- fast return to stable cold ground state



Research

#### 8. Dansgaard-Oeschger events. Statistical analysis

Calcium signal from GRIP: about 80 000 samples for 80 000 y



typical waiting time between D/O events: 1000 - 2000 y, waiting times between D/O events: multiples of  $\sim 1470$  years.

#### What triggers the transitions?

modeling by Langevin equation:

 $dX(t) = -U'(t, X(t))dt + \mathsf{NOISE}$ 

U — multi well potential, wells correspond to climate states

P. Ditlevsen (*Geophys. Res. Lett. 1999*): power spectrum analysis of time series:

NOISE contains strong  $\alpha$ -stable component with  $\alpha \approx 1.75$ .

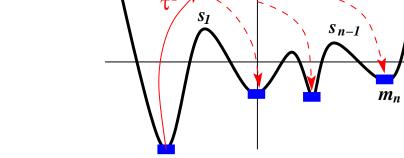
## 9. Simple system with Levy noise

consider SDE driven by  $\alpha$ -stable Lévy noise of small intensity  $X^{\varepsilon}(t) = x - \int_0^t U'(X^{\varepsilon}(s-)) \, ds + \varepsilon L(t), \quad \varepsilon \downarrow 0.$ 

• L is  $\alpha$ -stable symmetric Lévy process,  $\alpha \in (0,2)$ 

multi well potential U

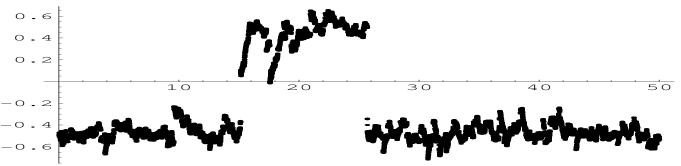
- $n \text{ local minima } m_i$
- n-1 local maxima  $s_i$
- $U''(m_i) > 0$ ,  $U''(s_i) < 0$



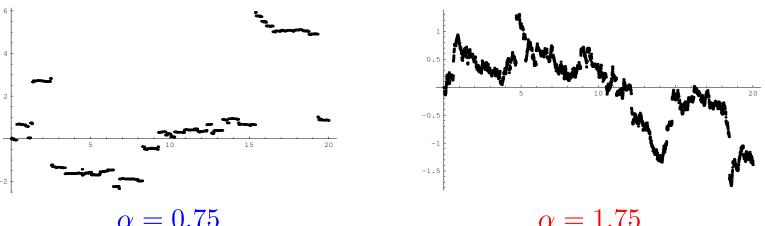
U(x)

x

aim: investigate exit and transition rates, meta-stability.



**10.** The symmetric  $\alpha$ -stable Lévy process L



 $\alpha = 1.75$ 

Lévy measure  $\nu(dy) = \frac{dy}{|y|^{1+\alpha}}$ , trajectories have countably many (small) jumps on finite time interval, jump times dense.

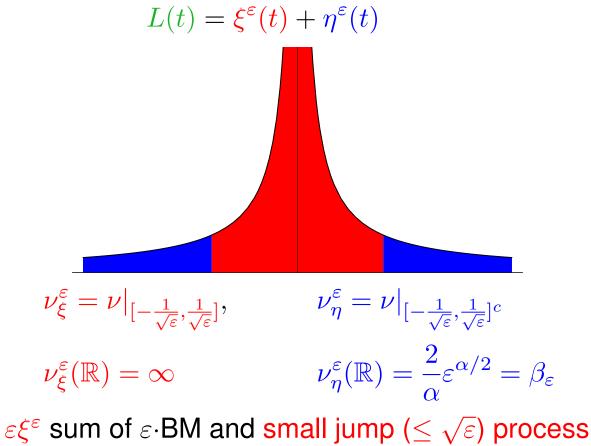
$$\mathbf{E}e^{i\lambda L(t)} = \exp\left\{-c(\alpha)|\lambda|^{\alpha}t\right\}, \quad c(\alpha) = 2\int_{0}^{\infty} \frac{1-\cos y}{y^{1+\alpha}} dy$$

$$\alpha = 1 \qquad \text{Cauchy process} \quad \frac{1}{\pi} \frac{1}{1+x^{2}}$$

$$\alpha = 2 \qquad \text{Brownian motion} \quad \frac{1}{\sqrt{2\pi}}e^{-\frac{x^{2}}{2}}$$

In physics: anomalous diffusion or Lévy flights

#### **11. Probabilistic approach of exit times**



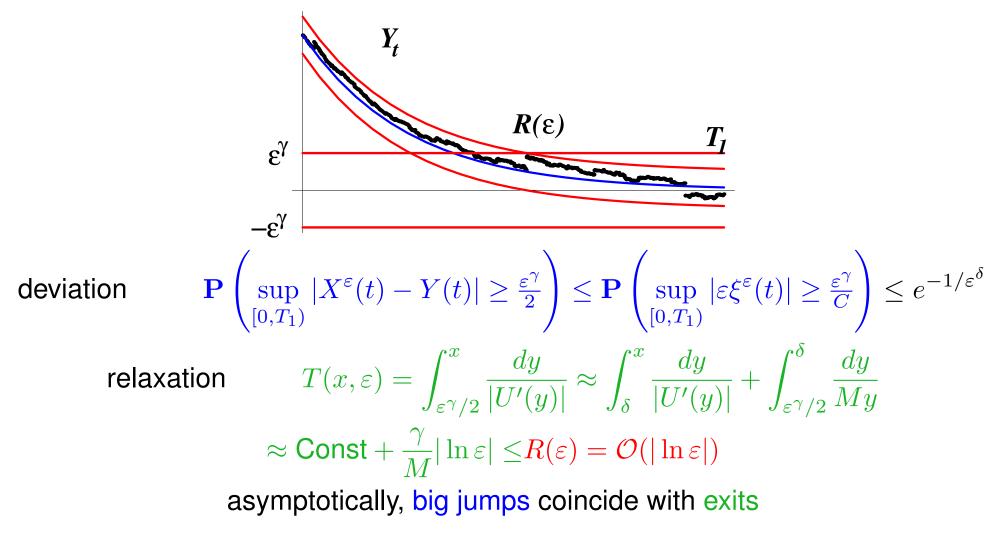
 $\varepsilon \eta^{\varepsilon}$  big jump ( $\geq \sqrt{\varepsilon}$ ) compound Poisson process big jumps at  $\tau_k$ , inter-jump time  $T_k$  with exponential law  $E(T_k) = (\beta^{\varepsilon})^{-1} = \frac{\alpha}{2} \varepsilon^{-\alpha/2}$ 

## 12. The small and large jump parts

U with stable state 0, exit from [-b, a] for a, b > 0

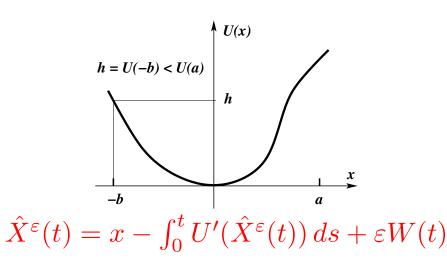
between big jumps  $X^{\varepsilon}$  is Y perturbed by  $\varepsilon \xi^{\varepsilon}$ 

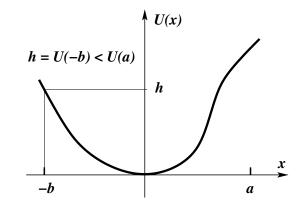
 $X^{\varepsilon}(t) = x - \int_0^t U'(X^{\varepsilon}(s-)) \, ds + \varepsilon \xi^{\varepsilon}(t), \quad t \in [0, T_1), \quad Y(t) = x - \int_0^t U'(Y(s)) \, ds$ 



# 13. Exit times: comparison of Gaussian and Lévy dynamics

 $\hat{\sigma} = \inf\{t \ge 0 : \hat{X}^{\varepsilon}(t) \notin [-b, a]\}$ 





 $\sigma = \inf\{t \ge 0 : X^{\varepsilon}(t) \notin [-b, a]\}$ 

$$X^{\varepsilon}(t) = x - \int_0^t U'(X^{\varepsilon}(s-)) \, ds + \varepsilon L(t)$$

Thm 1 (Freidlin-Wentzell):

$$\mathbf{P}_x(e^{(2h-\delta)/\varepsilon^2} < \hat{\sigma} < e^{(2h+\delta)/\varepsilon^2}) \to 1$$

Kramers' law ('40, Williams, Bovier et al.):

 $\mathbf{E}_x \hat{\sigma} \approx \frac{\varepsilon \sqrt{\pi}}{|U'(-b)| \sqrt{U''(0)}} e^{2h/\varepsilon^2}$ 

Exponential law (Day, Bovier et al.)

 $\mathbf{P}_x(\frac{\hat{\sigma}}{\mathbf{E}_x\hat{\sigma}} > u) \sim \exp\left(-u\right)$ 

Thm 2

$$\mathbf{P}_x(\frac{1}{\varepsilon^{\alpha-\delta}} < \sigma < \frac{1}{\varepsilon^{\alpha+\delta}}) \to 1$$

$$\mathbf{E}_x \sigma \approx \frac{1}{\varepsilon^{\alpha}} \left( \int_{\mathbb{R} \setminus [-b,a]} \frac{dy}{|y|^{1+\alpha}} \right)^{-1}$$

$$\mathbf{P}_x(\frac{\sigma}{\mathbf{E}_x\sigma} > u) \sim \exp\left(-u\right)$$

## **14.** *p*-Variation as test statistic

Ditlevsen's analysis: power spectrum of residua of time series Problem: Stationarity?

Aim: better test statistics than peaks of power spectrum.

Model assumption: with some U interpret data as

$$X^{\varepsilon}(t) = x - \int_0^t U'(X^{\varepsilon}(s-))ds + \varepsilon L(t) = Y^{\varepsilon}(t) + L^{\varepsilon}(t)$$

*L* Lévy process containing  $\alpha$ -stable component with unknown  $\alpha$ ,  $Y^{\varepsilon}$  of bounded variation; test  $\alpha$ 

Idea: *p*-variation characteristic for fluctuation behavior of noise processes.

$$V_t^{p,n}(X) = \sum_{i=1}^{[nt]} |X(\frac{i}{n}) - X(\frac{i-1}{n})|^p, \quad V_t^p = \lim_{n \to \infty} V_t^{p,n}$$

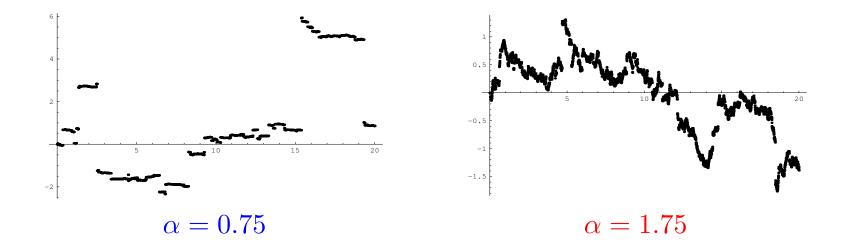
#### **15. Stable Processes and their** *p***-Variation**

*L* Lévy process with characteristics  $(d, \gamma, \nu)$  iff

$$E(\exp(iuL(t))) = \exp(t(-\frac{1}{2}du^2 + i\gamma u + \int_{\mathbf{R}} [e^{iuy} - 1 - iuy 1_{\{|y| \le 1\}}]\nu(dy))), \ u \in \mathbf{R}, t \ge 0,$$

 $\nu$  measure on Borel sets in **R** with  $\nu(\{0\}) = 0$ ,  $\int_{\mathbf{R}} [|y|^2 \wedge 1] \nu(dy) < \infty$ . *L*  $\alpha$ -stable symmetric Lévy process if

$$E(\exp(iuL(t))) = \exp(-c(\alpha)t|u|^{\alpha}), \quad \nu(dy) = \frac{1}{|y|^{\alpha+1}}dy, \quad u, y \in \mathbf{R}.$$



## **16.** *p*-Variation and the Blumenthal-Getoor Index

 $L \alpha$ -stable process with jump measure  $\nu$ ; then p-variation identified by Blumenthal-Getoor index

$$\beta_L = \inf\{s \ge 0 : \int_{\{|y| \le 1\}} |y|^s \nu(dy) < \infty\}$$

$$\gamma_L = \inf\{p > 0 : V_1^p(L) < \infty\}$$

#### Thm 1

L symmetric  $\alpha$ -stable. Then

$$\gamma_L = \beta_L = \alpha.$$

**Problem:** How to read  $\gamma_L = \alpha$  off the sequence  $(V_t^{p,n}(L))_{n \in \mathbb{N}}$ ?

Calls for results about the asymptotic behavior of the sequence.

# 17. The case $\alpha = 2$ : Brownian motion

For  $n \in \mathbb{N}$   $V_1^{p,n}(W)$  consists of n independent increments and

$$E(V_1^{p,n}(W)) = n^{1-\frac{p}{2}}E(|W(1)|^p)$$

Thm 2 (LLN type)

 $n^{-1+\frac{p}{2}}V_t^{p,n}(W) \rightarrow tE(|W(1)|^p)$  in probability, *Y* of bounded variation. Then also

 $n^{-1+\frac{p}{2}}V_t^{p,n}(W+Y) \rightarrow tE(|W(1)|^p)$  in probability.

Thm 3 (CLT type)

 $(n^{\frac{1}{2}}[n^{-1+\frac{p}{2}}V_t^{p,n}(W) - tE(|W(1)|^p)])_{t \ge 0} = (n^{-\frac{1}{2}+\frac{p}{2}}V_t^{p,n}(W) - n^{\frac{1}{2}}tE(|W(1)|^p))_{t \ge 0}$  $\rightarrow ((var(|W(1)|^p))^{\frac{1}{2}}\tilde{W}(t))_{t \ge 0}$ 

weakly with respect to the Skorokhod metric, and an independent Brownian motion  $\tilde{W}$ .

## 18. The case $\alpha < 2$

(Lit: Corcuera, Nualart, Wörner '07; case  $p < \alpha$  for LLN type,  $p < \frac{\alpha}{2}$  for CLT type)

**Problem:**  $p < \frac{\alpha}{2} < 1$  not satisfactory for paleo-climatic data! Beyond  $\frac{\alpha}{2}$  no CLT type result available, no asymptotic normality, but asymptotically of different type.

Thm 4 (LT type)  $L \alpha$ -stable with  $\alpha \in ]0, 2[$ . Then

$$(V_t^{p,n}(L) - B_t^n(\alpha, p))_{t \ge 0} \to \tilde{L}$$

weakly with respect to the Skorokhod metric, and an independent  $\frac{\alpha}{p}$ -stable process  $\tilde{L}$ . Here

$$B_t^n(\alpha, p) = \begin{cases} n^{1-\frac{p}{\alpha}} tE(|L(1)|^p), & \frac{\alpha}{2}$$

Same result with L + Y instead of L if Y is of finite p-variation and  $\frac{\alpha}{2} or <math>p > \alpha$ .

# **19. Methods of Proof**

- show that  $(|L(n) L(n-1)|^p)_{n \in \mathbb{N}}$  is in domain of attraction of  $\frac{\alpha}{p}$ -stable law; use tails and characteristic functions
- use Aldous' criterion for tightness of sequence  $(X^n)_{n \in \mathbb{N}}$  in Skorokhod metric:

(i) 
$$\lim_{K \to \infty} \sup_{n \in \mathbb{N}} P(\sup_{t \ge N} |X^n(t)| > K) = 0 \quad \text{for all } N \ge 0,$$

 $(ii) \lim_{\theta \downarrow 0} \limsup_{n \to \infty} \sup_{S \le T \le S + \theta} P(|X^n(T) - X^n(S)| \ge \varepsilon) = 0, \quad \text{for all } N \ge 0, \varepsilon > 0$ 

S,T stopping times

 adding processes of smaller variation: new notion of Lipschitz continuity on large sets; comparison of small and large jumps